

Variation on harmonic theme

Lukáš Vácha

ČVÚT, FJFI

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About me

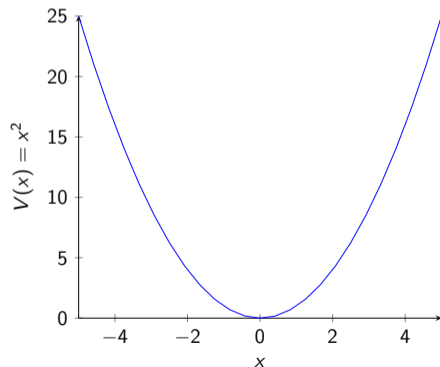
- 1 Lukáš Vácha
- 2 home: Hluboká nad Vltavou, South Bohemia



- 3 education: General grammar school in České Budějovice, now ČVÚT FJFI
- 4 work: waiter for 7 years, manager in Covid testing center
- 5 Bachelor thesis - Variation on harmonic theme

Linear harmonic oscillator

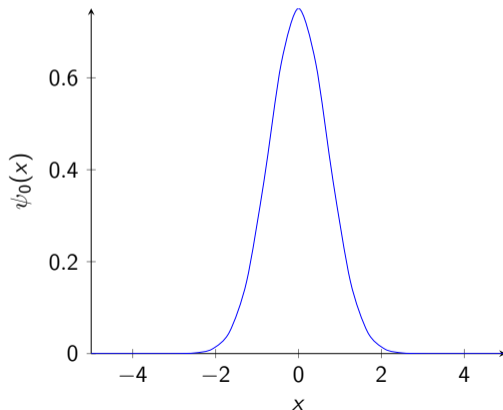
- 1 model describing body on spring or pendulum
- 2 Elastic force: $F = -kx$; Potential: $V(x) = \frac{1}{2}kx^2$
- 3 $m\ddot{x} + kx = 0$; $\ddot{x} + \omega^2x = 0$; $\omega = \sqrt{\frac{k}{m}}$
- 4 Solution: $x(t) = A\sin(\omega t + \varphi)$



- 1 Hilbert space: $\mathcal{H} = L^2(\mathbb{R}, dx)$
- 2 We set units, that $M = \hbar = 2$.
- 3 Same potential $V(x) = \omega^2 \hat{Q}^2$, now operator of multiplication on Hilbert space ($\hat{Q} = x \cdot$)
- 4 momentum p now also operator $\hat{P} = -2i \frac{\partial}{\partial x}$
- 5 Hamiltonian: $\hat{H}_{HO} = -\frac{d^2}{dx^2} + \omega^2 x^2$.
- 6 Hamiltonian is self-adjoint on $Dom(\hat{H}_{HO}) = \left\{ \psi \in L^2(\mathbb{R}, dx); \psi, \psi' \text{ a.c.}, \hat{H}\psi \in \mathcal{H} \right\}$
- 7 Spectrum of Hamiltonian and eigenvectors?

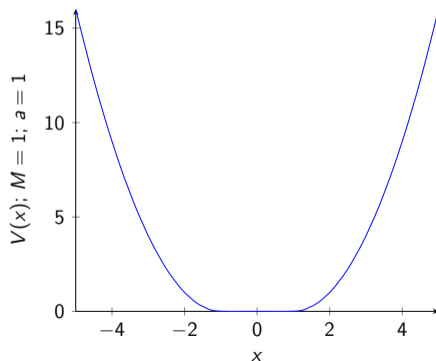
Quantum oscillator - solution

- 1 $\sigma(\hat{H}_{HO}) = \sigma_p(\hat{H}_{HO}) = \{E_n = 2\omega(n + \frac{1}{2}) \mid n \in \mathbb{N}_0\}$
- 2 Eigenvector in L^2 : $\psi_n(x) = A_n H_n(x) e^{-\frac{\omega x^2}{2}}$
- 3 Ground state: $E_0 = \omega$: $\psi_0(x) = \frac{1}{\sqrt[4]{\pi}} e^{-\omega x^2}$



Plateau modification

- 1 Modified potential:
$$V(x) = \begin{cases} \omega^2(x+a)^2; & x < -a \\ 0; & |x| < a \\ \omega^2(x-a)^2; & x > a \end{cases}$$
- 2 $\hat{H} = \frac{d^2}{dx^2} + V(x)$
- 3 \hat{H} is self-adjoint on the same domain as \hat{H}_{HO} [Kato-Rellich]



- 1 Now we deal with Schrödinger equation in 3 separate regions then connect (C^1) the solutions in points $-a$ and a

Quantization condition

$$\left[\frac{\omega}{\Gamma^2\left(\frac{1}{4} - \frac{E}{4\omega}\right)} - \frac{E}{4\Gamma^2\left(\frac{3}{4} - \frac{E}{4\omega}\right)} \right] \sin(2\sqrt{E}a) + \frac{\sqrt{E\omega}}{\Gamma\left(\frac{3}{4} - \frac{E}{4\omega}\right)\Gamma\left(\frac{1}{4} - \frac{E}{4\omega}\right)} \cos(2\sqrt{E}a) = 0$$

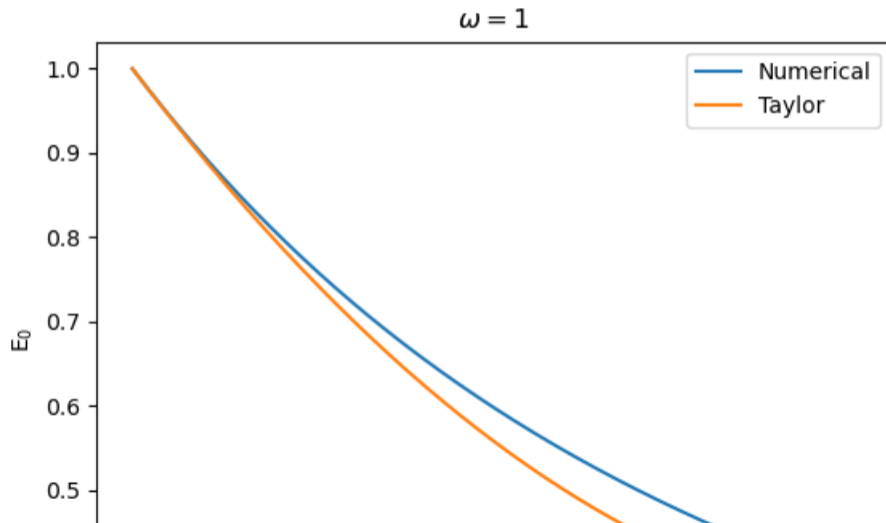
- 2 From this equation we get allowed energy levels, values from Hamiltonian spectrum.

- 1 Investigation of this model for small values of a
 - by expanding the equation to Taylor series
 - by quantum perturbation theory
- 2 Both methods lead to the same result

1st order perturbation theory for ground state energy

$$E_0(a) \sim \omega - \frac{2\omega\sqrt{\omega}}{\sqrt{\pi}} a$$

Numerical solution compared with perturbation theory



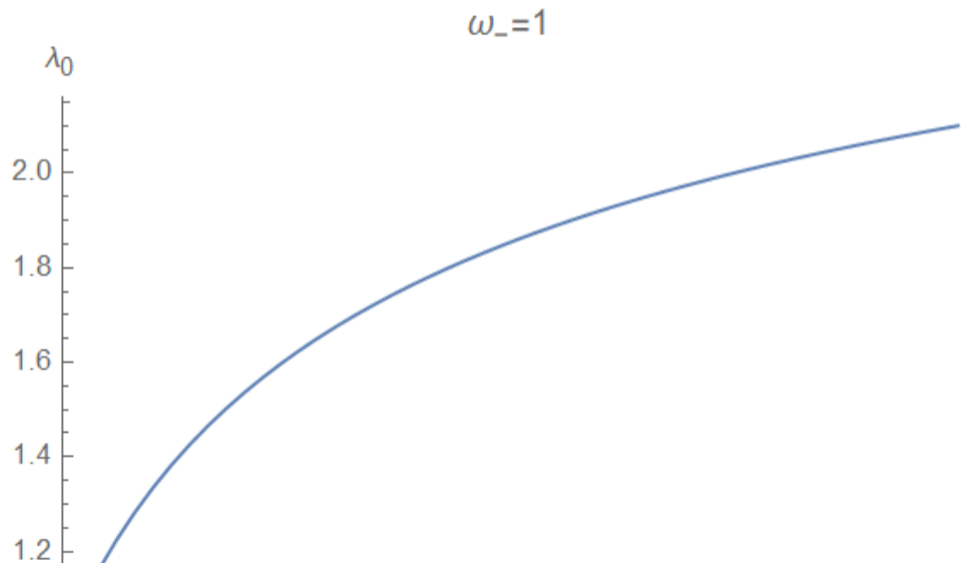
Different frequencies on half-axes modification

- 1 $\omega \neq \text{const.} \rightarrow \omega(x) = \begin{cases} \omega_+; x \geq 0 \\ \omega_-; x < 0 \end{cases}$
- 2 Potential: $V(x) = \begin{cases} \omega_-^2 x^2; x < 0 \\ \omega_+^2 x^2; x > 0 \end{cases}$

Hamiltonian spectrum equation

$$\frac{\sqrt{\omega_+}}{2\Gamma\left(\frac{3}{4} - \frac{E}{4\omega_-}\right)\Gamma\left(\frac{1}{4} - \frac{E}{4\omega_+}\right)} + \frac{\sqrt{\omega_-}}{2\Gamma\left(\frac{3}{4} - \frac{E}{4\omega_+}\right)\Gamma\left(\frac{1}{4} - \frac{E}{4\omega_-}\right)} = 0$$

Ground state energy depending on ω_+



- 1 Hamiltonian with periodic potential

Bloch wave

$$\psi_q(x) = e^{iqx} u_q(x)$$

- 2 Eigenfunction of Hamiltonian is a complex exponential multiple of periodic function u_q with same period as potential, q is index of such functions, referred to quasi-momentum
- 3 Monodromy matrix \mathbb{M} - period shift operator in base of linearly independent eigenvectors of Hamiltonian

Spectrum of Hamiltonian

$$\sigma(\hat{H}) = \left\{ E \mid \left| \frac{1}{2} \text{Tr} \mathbb{M} \right| \leq 1 \right\}$$

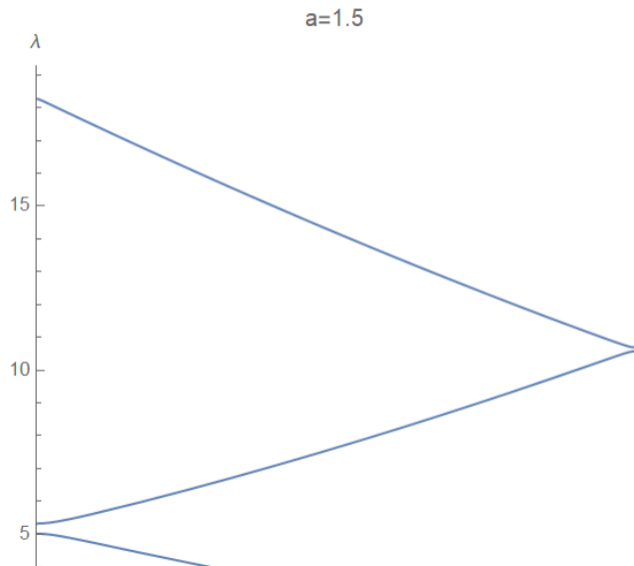
Periodic potential

- 1 Potential: $V(x) = \omega^2 x^2; |x| < a$
- 2 Periodically extended elsewhere
- 3 Bloch theory applied to this potential leads to condition

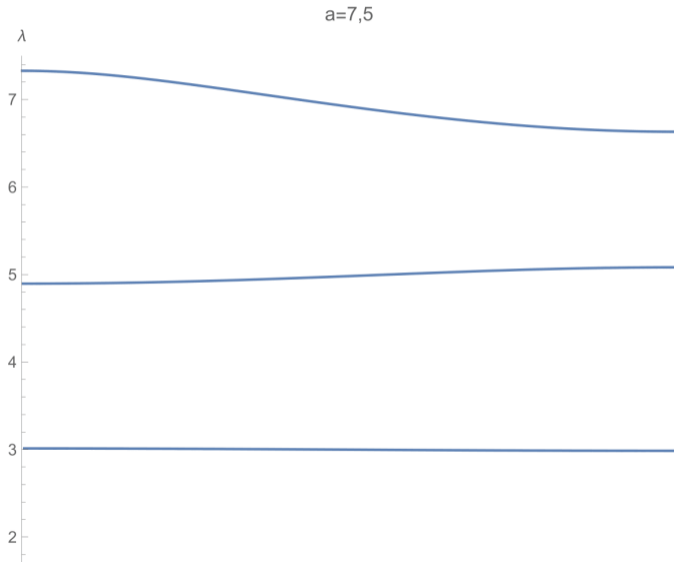
Spectral condition

$$\frac{a {}_1F_1\left(\frac{3}{4} - \frac{E}{4\omega}, \frac{3}{2}, \omega a^2\right) \left[-\omega {}_1F_1\left(\frac{1}{4} - \frac{E}{4\omega}, \frac{1}{2}, \omega a^2\right) + 2\omega \left(\frac{1}{2} - \frac{E}{2\omega}\right) {}_1F_1\left(\frac{5}{4} - \frac{E}{4\omega}, \frac{3}{2}, \omega a^2\right) \right]}{{}_1F_1\left(\frac{1}{4} - \frac{E}{4\omega}, \frac{1}{2}, \omega a^2\right) \left[(1 - \omega a) {}_1F_1\left(\frac{3}{4} - \frac{E}{4\omega}, \frac{3}{2}, \omega a^2\right) + a \left(\frac{1}{2} - \frac{E}{6\omega}\right) {}_1F_1\left(\frac{7}{4} - \frac{E}{4\omega}, \frac{5}{2}, \omega a^2\right) \right]} \leq 0$$

Energy bands for periodic potential



Energy bands for periodic potential



Conclusion

- 1 Results can be used in models of crystals in magnetic fields
- 2 Dielectric electrons in Lorenz interaction theory
- 3 Radiation in vacuum is described by set of linear harmonic oscillators

